

# Radial velocities vs Astrometry

- Consider an exoplanet with mass  $m_p$  orbiting a star of mass  $M_*$  at distance  $d$ , with semi-major axis  $a$ , period  $P$ , eccentricity  $e$ , inclination  $i$  with respect to the sky.
- Radial velocity wobble :

$$V_r = \underbrace{\frac{m_p}{m_p + M_*} \frac{2\pi a \sin i}{P \sqrt{1-e^2}} (\cos(\omega + \nu) + e \cos \omega)}_K \propto a^{-1/2}$$

$$K = \sqrt{\frac{2\pi G}{P} \frac{m_p \sin i}{(m_p + M_*)^{3/3}} \frac{1}{\sqrt{1-e^2}}} \approx 28.4 \text{ m s}^{-1} \times \left(\frac{P}{1 \text{ an}}\right)^{-1/3} \left(\frac{m_p \sin i}{M_{\text{Jup}}}\right) \left(\frac{M_*}{M_{\odot}}\right)^{-2/3}$$

- Astrometric displacement

$$\Delta \alpha = \frac{m_p}{m_p + M_*} \frac{a}{d} \frac{(1-e^2)}{1+e \cos \nu} \sqrt{1 - \sin^2(\omega + \nu) \sin^2 i} \sim \frac{m_p}{M_*} \frac{a}{d} \propto a$$

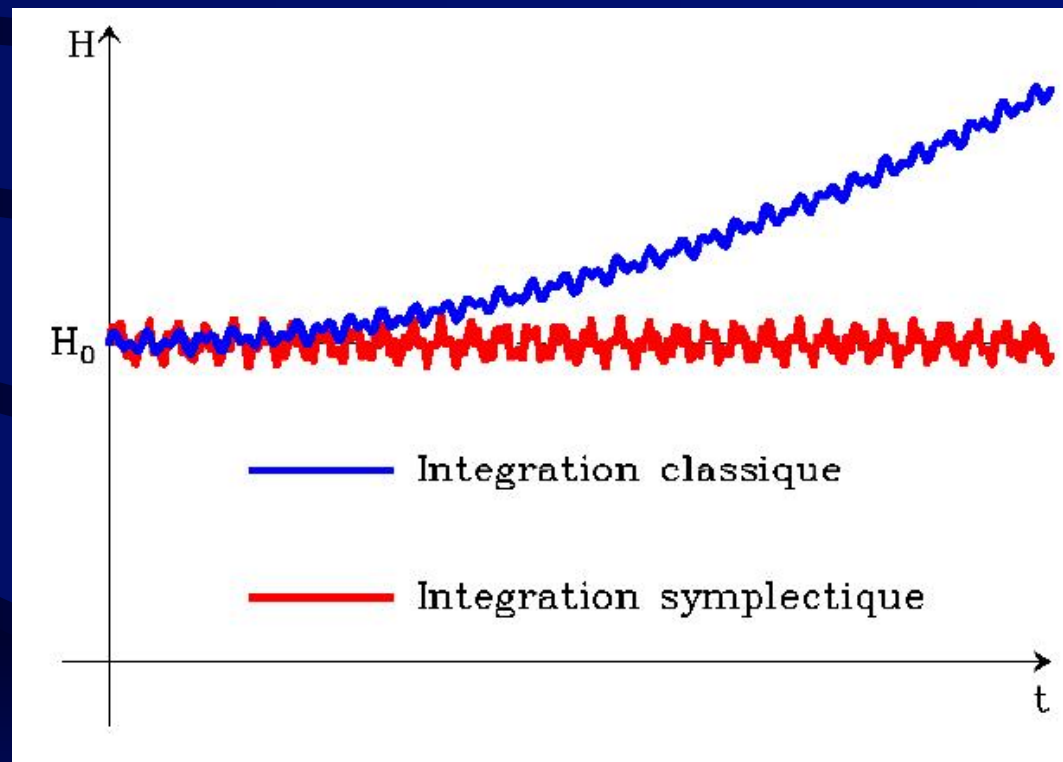
- Astrometry is **sensitive to large periods**, while velocimetry is sensitive to close-in planets.
- But it is best for **nearby stars**

# Observing nearby stars

- Requirement for astrometry : be patient ! (long periods)
- But for a nearby star, we can detect shorter period planets.
- Goal with such systems: Detecting solar system analogs.
- Issues to be investigated:
  - How frequent are solar-like systems around solar type stars ?
  - Do they all scale the same way ?
  - What about planets in binary systems ?
- Dynamical issues :
  - Orbital evolution in detected systems : stability, unseen planets ?
  - Perturbations by companions  $\Rightarrow$  Study with symplectic codes

# Symplectic N-body codes (1)

- Numerical integration = finding  $q(t)$  once  $q(t-\tau)$  is known
- Classical methods do not keep  $H = \text{Hamiltonian} = \text{cte}$  (Runge-Kutta, Burlish & Stoer ...)
- A symplectic scheme **exactly preserves H** or a nearby one
- The global error is **bound**
- Interesting for long term stability



# Symplectic N-body codes (2)

- Basic hypotheses of symplectic integration :

$H = H_A + H_B$  where we know how to exactly integrate  $H_A$  and  $H_B$ . We also try to have

- On suppose en plus  $H_B \ll H_A$  ( $H_B/H_A \sim \varepsilon$ )

- We have

$$\exp(\tau A) \cdot \exp(\tau B) = \exp[\tau(A+B) + O(\tau^2)]$$

$\Rightarrow$  Advancing  $H_A$  during  $\tau$  and then  $H_B$  during  $\tau$ , we build a **first order symplectic integrator** : we exactly solve a Hamiltonian

$$H_{\text{integ}} = H + O(\tau)$$

# Symplectic N-body codes (3)

- We also have

$$\exp\left(\frac{\tau}{2}B\right)\exp(\tau A)\exp\left(\frac{\tau}{2}B\right) = \exp\left[\tau(A+B) + O(\tau^3)\right]$$

⇒ We build a **second order symplectic integrator** if we integrate 1)  $H_A$  during  $\tau/2$ ,  $H_B$  during  $\tau$ , 3)  $H_A$  during  $\tau/2$ . In this case

$$H_{\text{integ}} = H + O(\tau^2)$$

- There are also 4, 6, 8... order methods
- In fact, if  $H_B/H_A \sim \varepsilon$ , we even have

$$H_{\text{integ}} = H + O(\varepsilon\tau^2)$$

⇒ We can integrate with a **large timestep**

# Practical N-body symplectic schemes

- **Simplest method** :  $T+U$ 
  - $H$ =Hamiltonian of  $N$  body problem = Kinetic + Potential =  $T+U$
  - Canonically conjugate variables :  $x_i=x,y,z..$   $p_i=mv_x,mv_y,mv_z...$
  - $T = T(p_i)$  and  $U=U(x_i) \Rightarrow$  We can integrate separately  $T$  and  $U$
  - $\Rightarrow$  Symplectic method with  $H_A=T, H_B=U$ . But we don't have  $H_B \ll H_A$
- **Second exemple**: MVS method (Mixed Variable Symplectic)
  - Hypothesis : ( $m_i \ll m_0$  pour every  $i = 1, \dots, n-1$ ) (planetary systems)
  - $H_A$  = Sum of Keplerian Hamiltonians (we can integrate...)
  - $H_B$  = What's left = Mutual perturbations  $\Rightarrow$  function of  $x_i$ 's only)
- The condition  $H_B \ll H_A$  is fulfilled if
  - 1) Body #0 is much more massive than the others
  - 2) Mutual distances are not too small (no **close encounters**)

# Close encounters ?

- Occurs when two bodies come close to each other. We have **strong and short scale perturbations**. Condition (2) is no longer fulfilled.
- The best to do : **reducing the timestep** when this occurs.
- But keeping a symplectic integrations requires a constant timestep ! Two ways to solve the problem:
  - **We drop symplecticity for a while** during the time of the encounter (**RMVS** code , *Levison & Duncan*). Ok from a statistical point of view.
  - **We change the  $H_A+H_B$  splitting** during the close encounter. We keep symplecticity but the computing time is longer (**MERCURY** code *Chambers*)
  - We split potentials in concentric spherical shells and attribute one timestep to each shell (**SyMBA** code *Duncan, Lee & Levison*).

## The HJS method (*Hierarchical Jacobi symplectic, Beust 2003*)

- The splitting of  $H$  is done according to the **hierarchical structure** of the stellar system.
- $H_A$  = Sum of Keplerian Hamiltonians corresponding to centers of mass orbiting other centers of mass
- $H_B$  = What's left
- The condition  $H_B \ll H_A$  is fulfilled irrespective of the various masses, once the nested orbits are **very different in size**.



# Analysis of the astrometric signal

- Issue : detecting an astrometric wobble of a few  $\mu\text{as}$
- Perturbation sources (for a star at 50 pc) :

Proper motion over 1 period	0.1 $\mu\text{as}$ – 1 mas
Non linear proper motion on sky	0.1 mas after 30 yrs
Parallactic orbit	20 mas
Effect of Earth eccentricity	0.3 mas
Barycentric motion of the Sun	0.15 mas
Planetary perturbations on the Earth	0.03 mas after 30 yrs

- $\Rightarrow$  Need for an accurate modeling of all these effects !